

# 08.128.809 Theoretische Elementarteilchenphysik

## Quantum Field Theory II

### Homework set 3

Due June 3, 2020, by 14:00

Please note how long it took you to solve each problem!

Monday, June 1, 2020 is a holiday. We will have the discussion session for HW 3 on Monday, June 8.

2-3, 30 pts. (repeat) Using the functional method, derive the Feynman rules in momentum space for the following interaction vertices:

A.  $\mathcal{L} = \lambda\phi^4$

B.  $\mathcal{L} = g^2\phi^* \phi A_\mu A^\mu$ . Note that for the purposes of functional methods, vector fields are simply collections of scalar fields indexed by an extra Lorentz index.

C.  $\mathcal{L} = g f_{abc} \partial_\mu A_\nu^a A^{\mu, b} A^{\nu, c}$ . For convenience, define all momenta to flow into the vertex, and assume  $f_{abc}$  is totally antisymmetric under interchange of any two [group space] indices.

D, Extra credit, 10 pts.  $\mathcal{L} = \frac{1}{4} g^2 (f^{eab} A_\mu^a A_\nu^b) (f^{ecd} A^{\mu, c} A^{\nu, d})$ . Again, assume  $f_{abc}$  is totally antisymmetric under interchange of any two [group space] indices.

3-1, 25 pts. Using the functional method, derive the Feynman rules in momentum space for the following interaction vertices:

A.  $\mathcal{L} = g \bar{\psi} \gamma_\mu A^\mu \psi$

B.  $\mathcal{L} = y \phi \bar{\psi} \psi + \text{h.c.}$

3-2, 75 pts. The concept of dimensional transmutation is one of the most important and involved topics in quantum field theory. We will mainly focus on the question of how a dimensionful scale arises in a physical theory when the fundamental bare Lagrangian only contains dimensionless couplings.

A, 20 pts. Read “Dimensional Analysis in Field Theory,” P. M. Stevenson. A PDF of the manuscript is available on Reader. In your own words, write a paragraph or two explaining how a dimensional scale can arise when the Lagrangian has no dimensionful couplings.

B, 25 pts. Read “Radiative Corrections as the Origin of Spontaneous Symmetry Breaking,” by S. Coleman and E. Weinberg. A PDF of the manuscript is available on Reader. In your own words, based on the discussion of massless scalar QED in section 4, write a paragraph or two explaining how a phase change in the theory can arise by a perturbative calculation.

C, 30 pts. The full calculation is done in sections II and III of the Coleman and Weinberg paper and as an exercise on pages 469-470 of Peskin and Schroeder. We will

not reproduce the full calculation but instead check the most important consequences. We start with the scalar QED Lagrangian:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_\mu\phi)^\dagger D^\mu\phi - m^2\phi^\dagger\phi - \frac{\lambda}{6}(\phi^\dagger\phi)^2, \quad (1)$$

with  $D_\mu = \partial_\mu + ieA_\mu$ .

- i, 15 pts. Expand the Lagrangian assuming the scalar field  $\phi$  obtains a vacuum expectation value (VEV)  $\langle\phi\rangle = \phi_0$ . The remaining scalar and pseudoscalar degrees of freedom of the complex scalar field  $\phi$  can be parametrized as  $\sigma$  and  $\pi$ . This result is encoded by plugging in

$$\phi = \phi_0 + \frac{1}{\sqrt{2}}(\sigma + i\pi) \quad (2)$$

into the Lagrangian. Setting  $m^2 = -\mu^2 < 0$ , solve for  $\phi_0$  in terms of  $\mu$  and  $\lambda$ . Identify the mass term for the  $A_\mu$  vector. This is the first example of the Higgs mechanism, where the nonzero VEV of the scalar field  $\phi$  leads to a spontaneous breaking of the QED gauge symmetry and a mass for the photon.

- ii, 15 pts. Given the form of the potential in equation 4.5 of Coleman and Weinberg,

$$V = \frac{\lambda}{4}\phi_c^4 + \left(\frac{5\lambda^2}{1152\pi^2} + \frac{3e^4}{64\pi^2}\right)\phi_c^4\left(\log\frac{\phi_c^2}{M^2} - \frac{25}{6}\right), \quad (3)$$

calculate the VEV of  $\phi_c$  by minimizing  $V$ . Note that the potential (which includes the relevant tree-level terms from the Lagrangian in equation 4.1) does not have a tree-level mass for  $\phi_c$ . To distinguish the behavior between when  $\langle\phi_c\rangle = 0$  and  $\langle\phi_c\rangle \neq 0$ , specify the necessary condition on the relative magnitude of  $\lambda$  and  $e$ .

- iii, Extra credit, 10 pts. Taking the  $\beta$ -functions for  $e$  and  $\lambda$  from Peskin and Schroeder,

$$\beta_e = \frac{e^3}{48\pi^2}, \quad \beta_\lambda = \frac{1}{24\pi^2}(5\lambda^2 - 18e^2\lambda + 54e^4), \quad (4)$$

sketch the renormalization group flows in the  $(\lambda, e^2)$  plane. You should find that every trajectory passes through the condition from part ii that marks when  $\phi_c$  acquires a vev.