

# 08.128.165 Theorie 6a, Relativistische Quantenfeldtheorie Quantum Field Theory I

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## Homework set 4

**Due November 23, 2022 by the start of lecture.**

**Please note how long it took you to solve each problem.**

- 4-1, 20 pts. Practice with Dirac algebra (part 2). Using the properties that  $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$  and  $\{\gamma^\mu, \gamma^5\} = 0$ , as well as the fact that the trace of any number of  $\gamma$  matrices obeys cyclicity,

$$\text{Tr}[\gamma^\mu \gamma^\nu \dots \gamma^\sigma] = \text{Tr}[\gamma^\nu \dots \gamma^\sigma \gamma^\mu] , \quad (1)$$

evaluate the following Dirac trace identities (the trace is performed over the spin indices).

- A, 4 pts.  $\text{Tr}[\gamma^\mu \gamma^\nu] = 4g^{\mu\nu}$  ,  
B, 8 pts.  $\text{Tr}[\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma] = 4(g^{\mu\nu} g^{\rho\sigma} - g^{\mu\rho} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\rho})$  ,  
C, 4 pts.  $\text{Tr}[\gamma^{\mu_1} \gamma^{\mu_2} \dots \gamma^{\mu_{2n-1}} \gamma^5] = 0$  , (trace of an odd number of matrices with  $\gamma^5$ )  
D, 4 pts.  $\text{Tr}[\gamma^{\mu_1} \gamma^{\mu_2} \dots \gamma^{\mu_{2n-1}}] = 0$  . (trace of an odd number of gamma matrices)

The two relations (C) and (D) are valid for any integer  $n \geq 1$ .

- 4-2, 10 pts. Practice with Dirac algebra (part 3). The charge conjugate matrix  $C$  has the properties  $C\gamma^\mu C^{-1} = (-\gamma^\mu)^T$  and  $C\gamma^5 C^{-1} = \gamma_5^T$ . Show that

- A, 5 pts.  $\text{Tr}[\gamma^\mu \gamma^\nu \gamma_5] = 0$  ,  
B, 5 pts.  $\text{Tr}[\gamma^{\mu_1} \gamma^{\mu_2} \dots \gamma^{\mu_{2n}}] = \text{Tr}[\gamma^{\mu_{2n}} \gamma^{\mu_{2n-1}} \dots \gamma^{\mu_1}]$  .

- 4-3, 35 pts. Any complex  $4 \times 4$  matrix  $M$  can be decomposed into Dirac basis, which are the 16 matrices  $\Gamma^A \in \{1_{4 \times 4}, i\gamma_5, \gamma^\mu, \gamma^\mu \gamma^5, \sigma^{\mu\nu}\}$ , with  $\mu, \nu = 0, 1, 2, 3$  and  $\mu < \nu$ . Thus we can write the expansion as

$$M = a1_{4 \times 4} + bi\gamma_5 + c_\mu \gamma^\mu + d_\mu \gamma^\mu \gamma_5 + \frac{1}{2} e_{\mu\nu} \sigma^{\mu\nu} , \quad (2)$$

with complex numbers  $a, \dots, e_{\mu\nu}$ . (We use a  $\frac{1}{2}$  factor on the last term since the sum is not restricted to  $\mu < \nu$ .)

A, 20 pts. Define an associate set of 16 basis matrices with lower indices by

$$\Gamma_B \in \{1_{4 \times 4}, -i\gamma_5, \gamma_\alpha, -\gamma_\alpha\gamma_5, \sigma_{\alpha\beta}\} , \quad (3)$$

for  $\alpha, \beta = 0, 1, 2, 3$  with  $\alpha < \beta$ . Knowing the traces of different combinations of matrices, show that  $\text{Tr}[\Gamma^A \Gamma_B] = 4\delta_B^A$  for any choices of  $A$  and  $B$ . Extract expressions for the coefficients  $a, \dots, e_{\mu\nu}$  in terms of traces of  $M$  with the matrices of the associated basis  $\Gamma_B$ . You can use the identity  $\text{Tr}[\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma_5] = -4i\epsilon^{\mu\nu\rho\sigma}$ .

B, 15 pts. Using the above result, decompose the product  $M = \gamma^\sigma \gamma^\rho \gamma^\delta$  into the Dirac basis.

4-4, 35 pts. Fierz identities. Let us consider products of bilinears,  $(\bar{u}_1 M u_2)(\bar{u}_3 N u_4)$ , where  $M$  and  $N$  are arbitrary  $4 \times 4$  matrices and  $u_i$  are arbitrary, 4-component Dirac spinors. We can rewrite the bilinear product

$$(\bar{u}_1 M u_2)(\bar{u}_3 N u_4) = \sum_{A,B} C_B^A (\bar{u}_1 \Gamma_A u_4)(\bar{u}_3 \Gamma^B u_2) \quad (4)$$

where  $\Gamma_A$  and  $\Gamma^B$  are the 16 matrices of the Dirac basis and the associated basis in 5-3A.

A, 15 pts. Using the completeness of the 16 basis matrices, show that the coefficients  $C_B^A$  are given by

$$C_B^A = \frac{1}{16} \text{Tr}[\Gamma^A M \Gamma_B N] . \quad (5)$$

**Hint:** You can use the identity:

$$1_{ij} 1_{kl} = \frac{1}{4} (\Gamma_A)_{il} (\Gamma^A)_{kj} = \frac{1}{4} (\Gamma^B)_{il} (\Gamma_B)_{kj} , \quad (6)$$

where  $i, j, k, l$  are 4-component spinor indices running from 1 to 4 and  $1 \equiv 1_{4 \times 4}$  is the  $4 \times 4$  identity matrix, so  $1_{ij} = \delta_{ij}$ . Given the above identity, you should calculate the quantity  $M_{ij} N_{kl}$  where the indices correspond to given entries in the matrices  $M$  and  $N$ . To get started, write  $M_{ij} N_{kl} = M_{ix} 1_{xj} 1_{ky} N_{yl}$  and use the identity above on the  $1_{xj} 1_{ky}$  terms. To complete the problem, you will have to decompose  $M$  and  $N$  as sums over the basis matrices. *Aside: the identity relation is the full “completeness” relation that is referenced above. In particular, the identity relation shows that the  $\Gamma^A$  and  $\Gamma_B$  matrices form a bilinear, orthogonal basis; the only complication compared to normal vector spaces is that we do not have an inner product that gives 1 or 0, but instead we have a matrix multiplication that gives  $4 \times 4$  matrices that are  $1_{4 \times 4}$  or 0. This is the same intuition from general  $SU(2)$  matrices and the Pauli matrices as a basis.*

B, 10 pts. Work out the Fierz identity for

$$(\bar{u}_1 \gamma^\mu (1 - \gamma_5) u_2)(\bar{u}_3 \gamma_\mu (1 - \gamma_5) u_4) . \quad (7)$$

C, 10 pts. Work out the Fierz identity for

$$(\bar{u}_1 \gamma^\mu (1 - \gamma_5) u_2)(\bar{u}_3 \gamma_\mu (1 + \gamma_5) u_4) . \quad (8)$$