

# 08.128.165 Theorie 6a, Relativistische Quantenfeldtheorie Quantum Field Theory I

Felix Yu and Prisco Lo Chiatto

## Homework set 2

**Due November 7, 2022 by start of discussion session.**

**Please note how long it took you to solve each problem.**

2-1, 40 pts. We start with the free field solution for the real Klein-Gordon field,

$$\phi(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_{\vec{p}}}} \left( a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} + a_{\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}} \right) , \quad (1)$$

and the equal-time commutation relations

$$[\phi(x), \phi(y)]|_{x^0=y^0} = 0 , \quad (2)$$

$$[\pi(x), \pi(y)]|_{x^0=y^0} = 0 , \quad (3)$$

$$[\phi(x), \pi(y)]|_{x^0=y^0} = i\delta^{(3)}(\vec{x} - \vec{y}) . \quad (4)$$

A, 15 pts. Explicitly verify that the Hamiltonian

$$H = \int d^3x \left[ \frac{1}{2}(\partial_t \phi)^2 + \frac{1}{2}(\vec{\nabla} \phi)^2 + \frac{1}{2}m^2 \phi^2 \right] \quad (5)$$

becomes

$$H = \int \frac{d^3p}{(2\pi)^3} E_{\vec{p}} \left( a_{\vec{p}}^\dagger a_{\vec{p}} + \frac{1}{2}(2\pi)^3 \delta^{(3)}(0) \right) . \quad (6)$$

B, 15 pts. Explicitly verify that the total momentum operator

$$\vec{P} = - \int d^3x \pi(\vec{x}) \vec{\nabla} \phi(\vec{x}) \quad (7)$$

can be reexpressed as

$$\vec{P} = \int \frac{d^3p}{(2\pi)^3} \vec{p} a_{\vec{p}}^\dagger a_{\vec{p}} . \quad (8)$$

C, 10 pts. Finally, write a paragraph explaining the significance of the  $\delta^{(3)}(0)$  term from part A, making analogy to the quantum mechanical harmonic oscillator as necessary. Does this term have physical consequences?

D, Extra credit, 10 pts. For completeness, we can adopt a regularization procedure for the infinite zero-point energy from part C by considering the energy density. (The simplest approach is to confine the spectrum to a finite box of volume  $V$ ; for any box size, the 3-dimensional delta function grows with size  $V$ , so dividing by  $V$  is equivalent to setting the 3D delta function to 1.) In this case, the zero point energy density is

$$\epsilon_0 = \frac{E_0}{V} = \int d^3p \frac{1}{2} E_{\vec{p}} . \quad (9)$$

Calculate the energy density over all momentum modes  $|\vec{p}|$ . You should get an *ultraviolet divergence* from the upper limits of integration of  $|\vec{p}|$  at  $\infty$ , corresponding to the assumption that the Hamiltonian is valid to arbitrarily high energy scales. To address the divergence, we can replace the upper limit of integration by  $\Lambda_{UV}$ : what is the new expression for the energy density? (In the end, to extract an observable quantity, we should remember that any scalar potential can include a finite constant contribution  $V(\phi) \supset V_0 + \frac{1}{2}m^2\phi^2 + \dots$ , which we need to keep for this discussion. The sum of the regularized zero-point energy density from the KG Hamiltonian and the  $V_0$  term then comprise the *dark energy* cosmological constant, which is measured to be  $\epsilon_0 \simeq (10^{-3} \text{ eV})^4$ .)

2-2, 60 pts. Consider a real function  $f(x)$ , which has a unique global minimum at  $x = 0$ . We define the integral

$$I(\alpha) = \int dx \exp \left\{ -\frac{1}{\alpha} f(x) \right\} . \quad (10)$$

A, 20 pts. Method of steepest descent. Perform a Taylor expansion of  $f(x)$  in the exponent and show that for  $\alpha \rightarrow 0$ , the integral has the asymptotic expansion

$$I(\alpha) = e^{-f_0/\alpha} \sqrt{\frac{2\pi\alpha}{f_0^{(2)}}} \left\{ 1 + \left( \frac{5}{24} \frac{(f_0^{(3)})^2}{(f_0^{(2)})^3} - \frac{3}{24} \frac{f_0^{(4)}}{(f_0^{(2)})^2} \right) \alpha + \mathcal{O}(\alpha^2) \right\} . \quad (11)$$

Here  $f_0 = f(0)$  and  $f_0^{(n)}$  is the  $n$ -th derivative evaluated at  $x = 0$ . You will find the Gaussian integral  $\int dx x^n e^{-ax^2}$  useful.

B, 20 pts. Method of stationary phase. Repeat the derivation in part A with the replacement of  $\alpha \rightarrow i\alpha$ . Can you justify why the expansions in part A and part B are both valid?

C, 20 pts. Causality violation in relativistic quantum mechanics. One of the main motivations for quantum field theory is that relativistic quantum mechanics allows for non-vanishing amplitudes for non-causal particle propagation. The amplitude for a free particle to propagate from  $\vec{x}_0$  to  $\vec{x}$  in quantum mechanics in  $U(t) =$

$\langle \vec{x} | e^{-iHt} | \vec{x}_0 \rangle$ . Using the relativistic expression for energy,  $H = E = \sqrt{|\vec{p}|^2 + m^2}$ , verify that

$$U(t) = \frac{1}{2\pi^2 |\vec{x} - \vec{x}_0|} \int_0^\infty dp \, p \sin(p|\vec{x} - \vec{x}_0|) e^{-it\sqrt{p^2+m^2}}, \quad (12)$$

with the usual  $p = |\vec{p}|$  for the magnitude of the momentum vector. Using the method of stationary phase, obtain the leading term in the asymptotic expansion of  $U(t)$  for spacelike separation of  $|\vec{x}|^2 \gg t^2$ , and explain how this term violates causality.